

## Exercise sheet 9

for the class on 25/06/20

### Excercise 1: Translation in momentum space

From the lecture notes, prove the following relation:

$$e^{ik\hat{y}} |p\rangle = |p + \hbar k\rangle \quad (1)$$

### Excercise 2: Limit temperature in Doppler cooling

In the lecture we derived a classical equation for the damping of the atomic motion in a 1D standing wave configuration

$$\dot{p} = -\frac{2\beta}{m}p \quad \Rightarrow \quad p(t) = p(0)e^{-\frac{2\beta}{m}t}. \quad (2)$$

However this simple model neglects fluctuations of the force which give rise to heating and thus to a non-zero final temperature. To account for this, the equation of motion for the momentum operator can be supplemented by

$$\dot{p} = -\frac{2\beta}{m}p + \zeta(t), \quad (3)$$

with  $\langle \zeta(t)\zeta(t') \rangle = D\delta(t - t')$ , where  $D = \hbar\omega_{\text{rec}}\gamma$  is the diffusion coefficient. By formally integrating Eq. (3), calculate the width of the momentum distribution  $\Delta p$  to infer the final temperature defined by  $\frac{1}{2}k_B T = \frac{\Delta p^2}{2m}$ . How does it change in three dimensions?

### Excercise 3: Quantum Brownian motion

The dynamics of a macroscopic oscillator (e.g. a vibrating membrane) of frequency  $\omega_m$  in thermal equilibrium with the environment (temperature  $T$ ) can be described by a set

of Langevin equations (in terms of dimensionless quadratures):

$$\dot{q} = \omega_m p, \quad (4a)$$

$$\dot{p} = -\gamma_m p - \omega_m q + \zeta(t). \quad (4b)$$

The model includes a damping term  $\gamma_m$  as well as a noise term responsible for thermalization. The correlations of the noise term in the time domain are

$$\langle \zeta(t)\zeta(t') \rangle = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{-i\omega(t-t')} \frac{\gamma_m \omega}{\omega_m} \left[ \coth\left(\frac{\hbar\omega}{2k_B T}\right) + 1 \right] = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{-i\omega(t-t')} S_{\text{th}}(\omega). \quad (5)$$

In the following we will assume that the oscillator is very good such that  $\gamma_m \ll \omega_m$  and that the temperature is high compared to the mechanical frequency:  $\hbar\omega_m \ll k_B T$ .

a) Show that under these approximations the decay rate and thermal occupancy follow from  $\gamma_m = \frac{S_{\text{th}}(\omega_m) - S_{\text{th}}(-\omega_m)}{2}$  and  $\bar{n} = \frac{S_{\text{th}}(-\omega_m)}{2\gamma_m}$  (Hint: Expand  $\coth(x)$  with  $x = \frac{\hbar\omega}{2k_B T} \ll 1$ ).

b) By going to Fourier space (Definition:  $A(\omega) = \frac{1}{\sqrt{2\pi}} \int dt e^{i\omega t} A(t)$ ), show that the position quadrature is related to the mechanical susceptibility  $\epsilon(\omega)$  according to  $q(\omega) = \epsilon(\omega)\zeta(\omega)$  with  $\epsilon(\omega) = \frac{\omega_m}{\omega_m^2 - \omega^2 - i\gamma_m\omega}$ . Using this result show that the variance of position (and momentum) quadrature is given by  $\langle q(t)^2 \rangle = \langle p(t)^2 \rangle = \bar{n} + \frac{1}{2}$ .